

RAIRO

MODÉLISATION MATHÉMATIQUE ET ANALYSE NUMÉRIQUE

PHILIPPE LE FLOCH

J. C. NEDELEC

Asymptotic time-behavior for weighted scalar conservation laws

RAIRO – Modélisation mathématique et analyse numérique,
tome 22, n° 3 (1988), p. 469-475.

http://www.numdam.org/item?id=M2AN_1988__22_3_469_0

© AFCET, 1988, tous droits réservés.

L'accès aux archives de la revue « RAIRO – Modélisation mathématique et analyse numérique » implique l'accord avec les conditions générales d'utilisation (<http://www.numdam.org/legal.php>). Toute utilisation commerciale ou impression systématique est constitutive d'une infraction pénale. Toute copie ou impression de ce fichier doit contenir la présente mention de copyright.

NUMDAM

Article numérisé dans le cadre du programme
Numérisation de documents anciens mathématiques
<http://www.numdam.org/>

ASYMPTOTIC TIME-BEHAVIOR FOR WEIGHTED SCALAR CONSERVATION LAWS (*)

Philippe LE FLOCH ⁽¹⁾, J. C. NEDELEC ⁽¹⁾

Communicated by C. BARDOS

Abstract. — For weighted scalar conservation laws introduced as model equations for gas flows in axisymmetric coordinates or in a nozzle, the asymptotic time-behavior of a entropy weak solution is obtained according to the behaviors of the weight and flux functions, thanks to the explicit formula previously derived in [10].

Résumé. — Considérant une loi de conservation scalaire avec poids qui modélise l'évolution d'un gaz en géométrie axisymétrique ou dans une tuyère, on utilise la formule explicite obtenue dans [10] pour préciser le comportement asymptotique de la solution faible entropique suivant le comportement de la fonction-flux et de la fonction-poids de l'équation.

1. INTRODUCTION : WEIGHTED SCALAR CONSERVATION LAWS

We are interested in *weighted scalar nonlinear hyperbolic conservation laws* in the half space

$$(1.1) \quad \frac{\partial}{\partial t} (r(x) u(x, t)) + \frac{\partial}{\partial x} (r(x) f(u(x, t))) = 0, \quad x > 0, \quad t > 0$$

with a \mathcal{C}^2 convex flux-function $f: \mathbb{R} \rightarrow \mathbb{R}$ and a \mathcal{C}^2 positive weight-function $r:]0, \infty[\rightarrow]0, \infty[$. Such equations are considered as model equations for gas dynamics in axisymmetric coordinates where the weight-function satisfies

$$(H.1) \quad r \in L^\infty(0, 1) \quad \text{and} \quad \frac{1}{r} \in L^\infty(1, \infty)$$

(for example take $r(x) = x^\alpha$, $\alpha = 0, 1, 2$), or as model equations for gas

(*) Received in July 1987.

⁽¹⁾ Centre de Mathématiques Appliquées, École Polytechnique, 91128 Palaiseau Cedex (France).

flows in a nozzle where there exists two positive constant r_- and r_+ such that

$$(H.2) \quad 0 < r_- \leq r(\cdot) \leq r_+ < +\infty .$$

For equations of the form (1.1), an uniqueness and existence result was proved in Le Floch-Nedelec [10] thanks to an explicit formula generalizing the one of Lax [7]. Moreover, a suitable formulation of boundary condition (at $x = 0$) was proposed. Refer also to previous works of Whitham [17] and Schonbek [15]. In this paper, we look for the asymptotic time-behavior of the solution of (1.1). We generalize well-known results of time-behavior for equations (1.1) without weight-function (that is $r(\cdot) \equiv 1$). The rate of convergence is specified according to the properties of the weight-function and the flux-function. For classical results on time-behavior of conservation laws, we refer to Lax [7], Dafermos [2], Conway [1], Liu-Pierre [13]... We present our result of time-behavior in the following section 2. Before, we detail some important properties of the mixed problem associated to the weighted scalar conservation law (1.1). The proofs can be found in [8]-[11].

We look for weak solutions (that is in the sense of distributions) of (1.1) satisfying an initial condition

$$(1.2) \quad u(x, 0) = u_0(x) , \quad x > 0$$

with an initial data $u_0 : \mathbb{R}_+ \rightarrow \mathbb{R}$. And for the sake of uniqueness, one adds an entropy condition (Lax [7], Oleinik [14], Kruskov [6])

$$(1.3) \quad u(x - 0, t) \geq u(x + 0, t) , \quad x > 0, t > 0 .$$

We impose a zero boundary condition at $x = 0$, since we are only interested in this paper in the rate of convergence and not in the nonlinear aspect of the boundary condition. Concerning boundary conditions, see [10] and also Le Floch [8], [9], Dubois-Le Floch [3, 4].

Our assumptions concerning the initial data are as follows

$$(H.3) \quad r \cdot u_0 \in L^1(\mathbb{R}_+) , \quad r \cdot f(u_0) \in L^\infty(\mathbb{R}_+) .$$

And, the flux-function satisfies for simplifications

$$(H.4) \quad f'' > 0 , \quad f(0) = f'(0) = 0 .$$

(for example, take $f(u) = \frac{u^2}{2}$). Then, the problem (1.1)-(1.3) admits one and only one solution $u(\cdot, \cdot)$ such that $r \cdot u$ belongs to $L^\infty(\mathbb{R}_t^+ ; L^1(\mathbb{R}_x^+))$ and $rf(u)$ to $L^\infty(\mathbb{R}^+ \times \mathbb{R}^+)$. It satisfies also the following *stability properties* for almost every $t \geq 0$: the L^∞ -stability

$$(1.4) \quad \|r(\cdot) f(u(\cdot, t))\|_{L^\infty(\mathbb{R}_x^+)} \leq \|r(\cdot) f(u_0(\cdot))\|_{L^\infty(\mathbb{R}_x^+)}$$

and the L^1 -semi-group property

$$(1.5) \quad \|(u(\cdot, t) - v(\cdot, t)) \cdot r(\cdot)\|_{L^1(\mathbb{R}_x^+)} \leq \| (u_0(\cdot) - v_0(\cdot)) \cdot r(\cdot) \|_{L^1(\mathbb{R}_x^+)}$$

for two solutions u and v corresponding to two initial data u_0 and v_0 respectively. The property (1.5) generalizes a previous result of Keyfitz [4].

Moreover, take $a = f'$ and let $f_+^{-1} : \mathbb{R}_+ \rightarrow \mathbb{R}_+$ and $f_-^{-1} : \mathbb{R}_- \rightarrow \mathbb{R}_-$ be the two inverse functions of the convex function f . For each $(x, t) \in \mathbb{R}_+ \times \mathbb{R}_+$, we can define — thanks to the hypotheses (H.1) or (H.2) (see [10]) — the function :

$$\mathbb{R}_+ \times \{-1, +1\} \ni (c, \varepsilon) \mapsto y(c, \varepsilon) \in \mathbb{R}^+$$

by the algebraic relation

$$(1.6a) \quad t = \int_{y(c, \varepsilon)}^x \frac{d\xi}{a\left(f_\varepsilon^{-1}\left(\frac{c}{r(\xi)}\right)\right)}$$

and the function $G : \mathbb{R}_+ \times \{-1, 1\} \ni (c, \varepsilon) \rightarrow G(c, \varepsilon)$ as follows :

$$(1.6b) \quad G(c, \varepsilon) = \int_0^{y(c, \varepsilon)} u_0(\xi) r(\xi) d\xi - c \cdot t + \int_{y(c, \varepsilon)}^x f_\varepsilon^{-1}\left(\frac{c}{r(\xi)}\right) r(\xi) d\xi.$$

Then the value $u(x, t)$ of the solution u of (1.1)-(1.3) is provided by the following *explicit formula*

$$(1.7) \quad u(x, t) = f_{\varepsilon(x, t)}^{-1}\left(\frac{c(x, t)}{r(x)}\right)$$

where $(c(x, t), \varepsilon(x, t))$ minimizes the function G .

2. ASYMPTOTIC TIME-BEHAVIOR

Using the explicit formula (1.6)-(1.7), we are able as in Lax [7] to get an uniform decay in power of t for the solution u of problem (1.1)-(1.3), according to the behavior of the function f . Assuming that

$$(2.1) \quad k_- \cdot |v|^{p-2} \leq f''(v) \leq k_+ \cdot |v|^{p-2}, \quad \forall v \in \mathbb{R}; p \geq 2, 0 < k_- \leq k_+$$

we have :

THEOREM Under hypotheses (H 1) or (H 2) and (H 3)-(H 4), the solution $u(.,.)$ of (1.1)-(1.3), when the time t tends to infinity, decreases with the following rate

$$(2.2) \quad |u(x, t)| \leq \left(\frac{kM}{r(x)} \right)^{\frac{1}{p}} \cdot \frac{1}{t^{1/p}} \quad t > 0, x > 0$$

where the constants k and M depend only on the flux-function and the initial data respectively

$$k = 2p \frac{k_+}{k_-^2}$$

and

$$M = \|r \cdot u_0\|_{L^1(\mathbb{R}_+)} \quad \blacksquare$$

In the case of flows in axisymmetric coordinates

$$r(x) = x^\alpha, \quad \alpha \geq 0$$

the estimation (2.2) becomes

$$(2.3) \quad |u(x, t)| \leq \frac{c}{x^{\alpha/p}} \cdot \frac{1}{t^{1/p}} \quad (c > 0)$$

Note that, when $\alpha > 0$, this inequality (2.3) provides also the behavior of u when $x \rightarrow 0+$ and $x \rightarrow +\infty$

Proof of theorem Multiplying (2.1) by $v \in \mathbb{R}$ and integrating lead to

$$(2.4) \quad k_- \cdot \frac{|v|^p}{p} \leq a(v)v - f(v) \leq k_+ \cdot \frac{|v|^p}{p}$$

because of

$$\frac{d}{dv} (a(v)v - f(v)) = a'(v)v = f'' * (v)v$$

Moreover, by two successive integrations of (2.1), we have also

$$(2.5) \quad k_- \cdot \frac{|v|^p}{p(p-1)} \leq f(v) \leq k_+ \cdot \frac{|v|^p}{p(p-1)}$$

Then from (2.4)-(2.5) it results that

$$\frac{k_-}{k_+} (p-1) f(v) \leq a(v)v - f(v) \leq \frac{k_+}{k_-} (p-1) f(v)$$

or, for $v > 0$:

$$(2.6a) \quad \frac{k_-}{k_+} (p - 1) \frac{f(v)}{a(v)} \leq v - \frac{f(v)}{a(v)} \leq \frac{k_+}{k_-} (p - 1) \frac{f(v)}{a(v)}$$

and for $v < 0$:

$$(2.6b) \quad \frac{k_+}{k_-} (p - 1) \frac{f(v)}{a(v)} \leq v - \frac{f(v)}{a(v)} \leq \frac{k_-}{k_+} (p - 1) \frac{f(v)}{a(v)}.$$

Henceforth, the following minoration of the function G defined by (1.6) is an immediate consequence of (2.6) (used with the value $v = f_\varepsilon^{-1}\left(\frac{c}{r(\xi)}\right)$)

$$\begin{aligned} G(c, \varepsilon) &= \int_0^{y(c, \varepsilon)} u_0 r \, d\xi + \int_{y(c, \varepsilon)}^x \left\{ f_\varepsilon^{-1}\left(\frac{c}{r(\xi)}\right) - \frac{\frac{c}{r(\xi)}}{a f_\varepsilon^{-1}\left(\frac{c}{r(\xi)}\right)} \right\} r(\xi) \, d\xi \\ &\geq -M + \int_{y(c, \varepsilon)}^x \frac{k_-}{k_+} (p - 1) \frac{\frac{c}{r(\xi)}}{a f_\varepsilon^{-1}\left(\frac{c}{r(\xi)}\right)} \cdot r(\xi) \, d\xi \end{aligned}$$

where $M = \|r \cdot u_0\|_{L^1(\mathbb{R}^+)}$. Thus, the inequality

$$(2.7) \quad G(c, \varepsilon) \geq -M + \frac{k_-}{k_+} (p - 1) c \cdot t$$

holds for all $c = 0, \varepsilon = \pm 1, t > 0$.

Furthermore, for each (x, t) in $\mathbb{R}_+ \times \mathbb{R}_+$, the minimum value of the function G — which is by notation attained at $(c(x, t), \varepsilon(x, t))$ — may be majorized by the value of G at $c = 0$:

$$(2.8) \quad G(c(x, t), \varepsilon(x, t)) \leq \int_0^x r u_0 \, d\xi \leq M.$$

Now comparing (2.7) and (2.8), it results the estimation

$$(p - 1) t \frac{k_-}{k_+} c(x, t) \leq 2M$$

or

$$0 \leq c(x, t) \leq \frac{k_+}{k_-} \frac{2M}{(p - 1) t}.$$

Thus in virtue of (1.7), we have proved that $rf(u)$ decreases uniformly in $x \in \mathbb{R}_+$ when $t \rightarrow \infty$:

$$\sup_{x \in \mathbb{R}_+} |r(x) f(u(x, t))| \leq \frac{k_+}{k_-} \frac{2M}{(p-1)t}.$$

It remains to use again (2.5) :

$$k_- \cdot \frac{|u(x, t)|^p}{p(p-1)} \leq f(u(x, t)) \leq \frac{2Mk_+}{k_- \cdot r(x)} \cdot \frac{1}{t}$$

which gives (2.2). ■

When the solution u is bounded — say the weight-function r satisfies the properties (H.2) — it suffices to assume that the inequalities (2.1) hold in the neighborhood of $v = 0$ with one $p_0 \geq 2$ and also that f'' is uniformly bounded :

$$\text{Cte} \leq f'' \leq \text{Cte}'.$$

Namely, thanks to the theorem applied with $p = 2$, we know that the solution u tends to zero uniformly in $x \in \mathbb{R}_+$ when the time goes to infinity

$$\sup_{x \in \mathbb{R}_+} |u(x, t)| \leq \left(\frac{kM}{r_-} \right)^{\frac{1}{2}} \cdot \frac{1}{\sqrt{t}} \quad t > 0.$$

So for t sufficiently large, say $t > T_0$, $u(x, t)$ is for each x in \mathbb{R}_+ in the neighborhood where (2.1) holds. And the same proof as previously gives

$$|u(x, t)| \leq \left(\frac{kM}{r_-} \right)^{\frac{1}{p_0}} \cdot \frac{1}{t^{1/p_0}}, \quad \forall x \in \mathbb{R}_+, \forall t > T_0.$$

REFERENCES

- [1] E. D. CONWAY, *The formation and decay of shocks for a conservation law in several dimensions*, Arch. Rat. M.A. 64 (1977) pp. 47-57.
- [2] C. M. DAFERMOS, *Characteristics in hyperbolic conservation law*, in « Nonlinear analysis and mechanics : Heriot-Watt Symposium vol. 1 », Knops Editor (1983).
- [3] F. DUBOIS, Ph. LE FLOCH, *Boundary conditions for nonlinear hyperbolic systems of conservation laws*, Internal Report (1987), École Polytechnique ; J. of Diff. Eq., Vol. 71, No 1, jan. 1988, pp. 93-122.

- [4] F. DUBOIS, Ph. LE FLOCH, *Condition à la limite pour un système de lois de conservation*, Note Compt. Rend. Acad. Sc. Paris, t. 304, Série I, n° 3, pp. 75-78 (1987).
- [5] B. KEYFITZ, *Solutions with shocks, an example of L^1 -contractive semi-group*, Comm. Pure Appl. Math., 24 (1971) pp. 125-132.
- [6] S. N. KRUSKOV, *First order quasi-linear equations in several independant variables*, Math. USSR Sb., 10 (1970) n° 2, pp. 217-243.
- [7] P. D. LAX, *Conservation laws and the mathematical theory of shock waves*, CBMS Ser. Appl. Math., vol. 11, SIAM, Philadelphia (1973).
- [8] Ph. LE FLOCH, *Explicit formula for scalar conservation laws with boundary conditions*, to appear in Math. Meth. in Appl. Sc. (1988) vol. 10.
- [9] Ph. LE FLOCH, *Generalized Riemann problem and boundary conditions for systems of conservation laws*, Thesis (1987) École Polytechnique (France).
- [10] Ph. LE FLOCH, J. C. NEDELEC, *Explicit formula for weighted scalar conservation laws*, Internal Report n° 144 (janv. 1986) of École Polytechnique ; accepted for publication to Transactions of A.M.S.
- [11] Ph. LE FLOCH, J. C. NEDELEC, *Lois de conservation scalaires avec poids*, Note Compt. Rend. Acad. Sc. Paris, t. 301, Série I, n° 17, pp. 1301-1304 (1985).
- [12] Ph. LE FLOCH, P. A. RAVIART, *Un développement asymptotique pour le problème de Riemann généralisé*, Compt. Rend. Acad. Sc. Paris, t. 304, Série I, n° 4, pp. 119-122 (1987) and Ann. Inst. Henri Poincaré, Analyse non linéaire.
- [13] T. P. LIU, M. PIERRE, *Source-solutions and asymptotic behavior in conservation laws*, J. of Diff. Eq. 51, 419-441 (1984).
- [14] O. A. OLEINIK, *Discontinuous solutions of nonlinear differential equations*, A.M.S. Transl., Ser. 2, 26, pp. 95-172 (1963).
- [15] M. E. SCHONBEK, *Existence of solutions to singular conservation laws*, Siam J. Math. Anal., vol. 15, n° 6 (nov. 1984).
- [16] J. A. SMOLLER, *Reaction-Diffusion Equations and Shock Waves*, Springer, Verlag 258 (1983).
- [17] G. B. WHITHAM, *Linear and Non linear Waves*, Wiley Interscience, New York (1974).